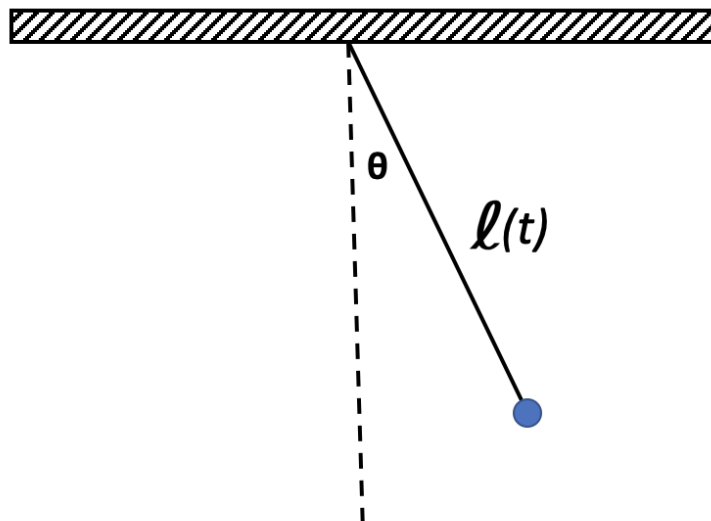


Winter 2022
Written Graduate Comprehensive Exam in Physics
UCLA
February 19, 2022

4 hours
One day
Six problems
Closed Book. Closed Notes.

Write Student Code Number: _____

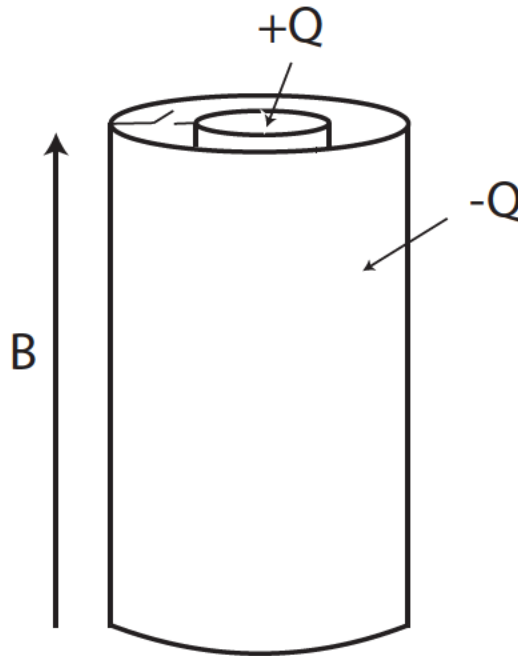
1. Consider a pendulum consisting of a point mass m attached to a string of slowly increasing length $\ell(t)$. The change of length versus time is $\dot{\ell}$. The motion is confined to a plane and we assume that $|\ell/\dot{\ell}|$ is much greater than the period of the oscillation.



- (a) (6 pts) Find the Lagrangian $L(\theta, \dot{\theta}, t)$ and the Hamiltonian $H(\theta, p_\theta, t)$ of the system, where θ is the angle of the string relative to downward, p_θ is the conjugate momentum, and $\dot{\theta}$ is the change of θ versus time.
- (b) (4 pts) Show why H is or is not equal to the total energy E of the pendulum.
- (c) (4 pts) Derive the equation of motion for θ in the form of a differential equation. When the change of length versus time $\dot{\ell} = 0$, what is the angular frequency ω_0 of small oscillations?
- (d) (6 pts) Suppose the length varies slowly so that $|\dot{\ell}/\ell| \ll \omega_0$, the angular frequency of oscillation. Show that as the length varies, the amplitude of small oscillations varies proportionally to ℓ^a and find a , the numerical exponent. (Hint: Consider using $\int p_\theta d\theta$ as an adiabatic invariant.)

2. Consider a cylindrical capacitor of length L and moment of inertia I with charge $+Q$ on the inner cylindrical metallic shell of radius a and charge $-Q$ on the outer cylindrical metallic shell of radius b . The capacitor is immersed in a uniform magnetic field \vec{B} , which points along the symmetry axis of the capacitor. There is a vacuum between the plates, which are held fixed to one another (e.g., by small insulating posts so that they will rotate together as a rigid body). A thin stainless steel wire with a switch is attached across the two plates of the capacitor. At $t = 0$, the switch is closed, and current flows from one plate to the other through a resistive wire (you can assume resistance R). This current creates a torque on the capacitor and it rotates.

Once the capacitor is fully discharged (the charge on each plate is zero), what is the magnitude and direction of the angular velocity, $\vec{\omega}$, of the capacitor? (For simplicity, ignore fringing fields and assume that the thin stainless wire does not break the symmetry of the system when computing fields.)

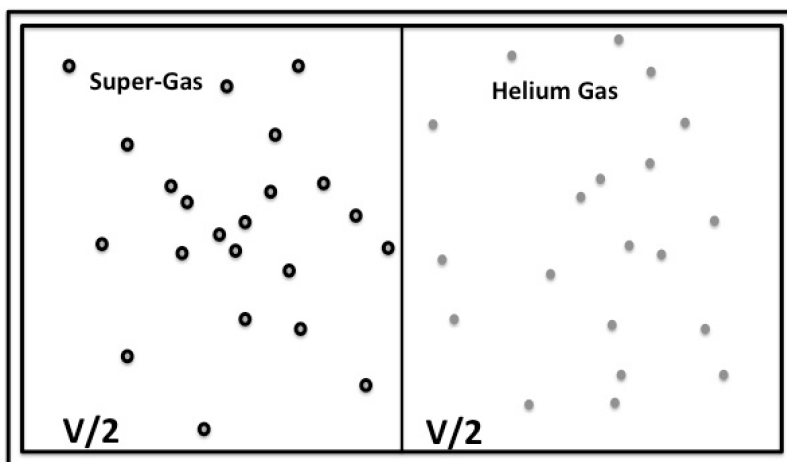


3. Physicists at a top-secret government laboratory discovered a Super-Gas that has the following number of microstates:

$$\Omega = \frac{\kappa^{10N} \pi^{10N} V^{10N} U^{10N}}{N! h^{20N} c^{10N}}$$

as a function of internal energy U , volume V , and the number of particles N . The constants π , h , and c are pi, Planck's constant, and the speed of light. The constant κ makes the units come out correctly.

The Super-Gas is placed on the left-hand side of the container below and helium gas is placed on the right-hand side. Both have the same number of particles N and occupy the same volume $V/2$. (For simplicity, neglect the heat capacity of the container and assume the container is isolated from the rest of the world. Use a simple model for the helium gas.)



a) The wall between the gas and the Super-Gas conducts heat but cannot move. At thermal equilibrium what is the ratio of the internal energy of the Super-Gas to that of the helium gas?

b) Now both sides are at temperature T . The wall is allowed to slide so that the volumes can change slowly. What is the ratio of the final volume of Super-Gas to that of helium?

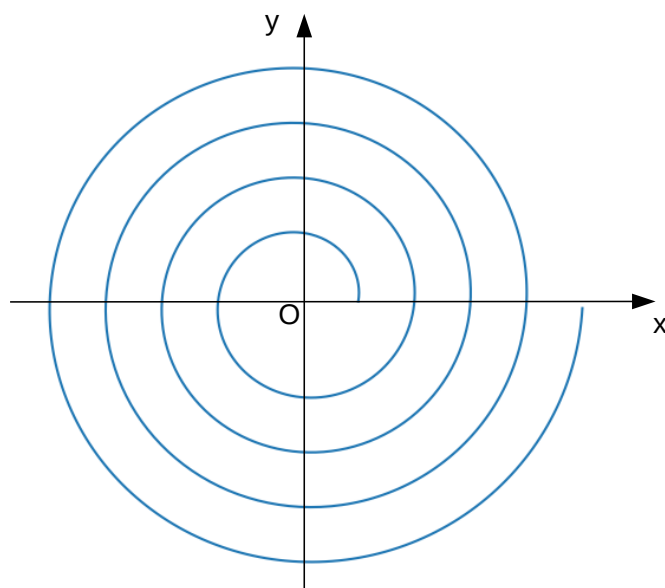
4. Part of a circuit consists of an Archimedean spiral which is given by the following equation in cylindrical coordinates:

$$\rho = a\phi,$$

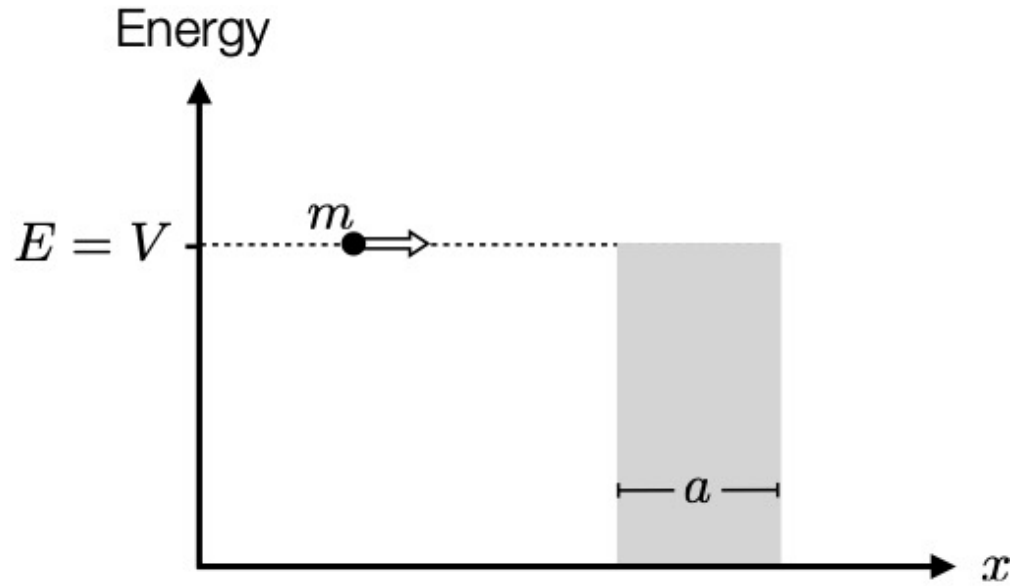
where $a > 0$, N is a positive integer > 1 , and $2\pi < \phi < 2\pi N$ as shown below for $N = 5$. The origin is marked by O.

A steady current I is running counter-clockwise in the spiral. Find the contribution to the magnetic field (direction and magnitude) from the spiral at the origin as a function of N , a , I , and physical constants.

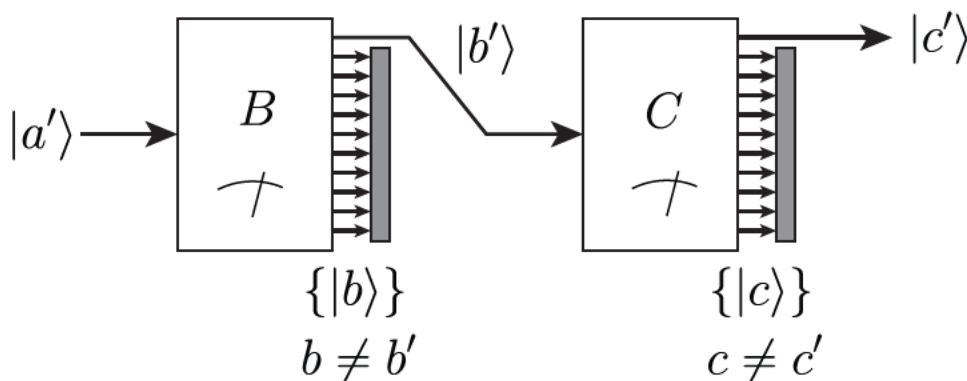
For simplicity, ignore the contribution from the rest of the circuit and assume no charge is building up at the ends.



5. Starting with the time-independent Schrödinger equation, work out the fraction of incident particles transmitted through a rectangular one-dimensional potential barrier in the case shown below, where the energy E of the incident particles is equal to the barrier height V . Let the particles have mass m and let the barrier width be a .



6. Consider a set of quantum operators A , B , and C that are mutually incompatible, which is to say $[A, B] \neq 0$, $[B, C] \neq 0$, and $[A, C] \neq 0$. Eigenkets and eigenvalues (assumed to be non-degenerate) of these operators will be denoted by lower-case letters for the upper-case operators (as in $B|b\rangle = b|b\rangle$, etc.), and we will consider the sequential measurement of an eigenstate of the A operator, $|a'\rangle$, in bases of B and C , as shown below.



As shown above, an input state $|a'\rangle$ is measured in the B basis and sent to the next measurement only if the eigenvalue of the measurement of B is b' , followed by a similar process in the C basis. All other measurement outcomes result in the state being filtered out.

a) (4 pts) Write expressions for the expansion coefficients c_b for the state $|a'\rangle$ in the B basis,

$$|a'\rangle = \sum_b c_b |b\rangle$$

and c_c for the state $|b'\rangle$ in the C basis,

$$|b'\rangle = \sum_c c_c |c\rangle.$$

b) (4 pts) In terms of these coefficients, what is the probability that in the experimental setup above the input state ($|a'\rangle$) results in final measurement outcome $|c'\rangle$?

c) (4 pts) What is the sum of the probabilities of outcome $|c'\rangle$ if you add up contributions from the complete set of choices for the intermediate eigenket $|b'\rangle$?

d) (8 pts) If we were instead to bypass the intermediate measurement B and feed $|a'\rangle$ directly into the C measurement, would the probability of obtaining outcome $|c'\rangle$ have to be the same as the answer to part (c)? Briefly explain or show why or why not.

This problem was from the 2017 comp exam.

Solution to #1

a) $L = T - V$

$$T = \frac{1}{2} m \dot{l}^2 + \frac{1}{2} I \dot{\omega}^2$$
$$= \frac{1}{2} m \dot{l}^2 + \frac{1}{2} m l^2 \dot{\theta}^2$$

$$V = mgh = mgl(1 - \cos\theta)$$

$$= mgl - mgl \cos\theta$$

ignore overall constant



$$h = l - l \cos\theta$$

$$L = \frac{1}{2} m \dot{l}^2 + \frac{1}{2} m l^2 \dot{\theta}^2 + mgl \cos\theta$$

$$H = p_\theta \dot{\theta} - L$$

$$\left(p_\theta = \frac{\partial L}{\partial \dot{\theta}} = m l^2 \dot{\theta} \Rightarrow \dot{\theta} = \frac{p_\theta}{m l^2} \right)$$

$$= m l^2 \dot{\theta}^2 - \frac{1}{2} m (\dot{l}^2 + l^2 \dot{\theta}^2) - mgl \cos\theta$$

$$= \frac{1}{2} m (l^2 \dot{\theta}^2 - \dot{l}^2) - mgl \cos\theta$$

$$= \frac{1}{2} m \left(l^2 \frac{p_\theta^2}{m^2 l^4} - \dot{l}^2 \right) - mgl \cos\theta$$

$$= \frac{1}{2} m \left(\frac{p_\theta^2}{m^2 l^2} - \dot{l}^2 \right) - mgl \cos\theta$$

$$H = \frac{1}{2} \frac{p_\theta^2}{m l^2} - \frac{1}{2} m \dot{l}^2 - mgl \cos\theta$$

$$\begin{aligned}
 b) \quad E &= T + V \\
 &= \frac{1}{2} m \dot{l}^2 + \frac{1}{2} m l^2 \dot{\theta}^2 - mgl \cos \theta \\
 &= \frac{1}{2} m \dot{l}^2 + \frac{1}{2} m l^2 \left(\frac{P_\theta^2}{m^2 l^4} \right) - mgl \cos \theta \\
 E &= \frac{P_\theta^2}{2 m l^2} + \frac{1}{2} m \dot{l}^2 - mgl \cos \theta
 \end{aligned}$$

$$H = \frac{P_\theta^2}{2 m l^2} - \frac{1}{2} m \dot{l}^2 - mgl \cos \theta$$

$$H \neq E \text{ (unless } \dot{l} = 0 \text{)}$$

c) Lagrange's equation:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} - \frac{\partial L}{\partial \theta} = 0$$

$$\frac{d}{dt} (m l^2 \dot{\theta}) + mgl \sin \theta = 0$$

$$m l^2 \ddot{\theta} + 2 m l \dot{l} \dot{\theta} + mgl \sin \theta = 0$$

$$\boxed{\ddot{\theta} + 2 \frac{\dot{l}}{l} \dot{\theta} + \frac{g}{l} \sin \theta = 0}$$

For small oscillations and $\dot{l} = 0$

$$\ddot{\theta} + \frac{g}{l} \theta = 0 \quad \text{physics} \Rightarrow \omega_0 = \sqrt{g/l}$$

d) $\mathcal{I} = \int_{\text{cycle}} d\theta p_\theta$ is adiabatic invariant = const.
For small θ and l small, trial
solution is

$$\theta(t) = A(t) \cos(\omega_0 t)$$

$$\mathcal{I} = \int_{\text{cycle}} d\theta (ml^2 \dot{\theta}) = \text{const}$$
$$\theta = \dot{\theta} t$$
$$d\theta = \dot{\theta} dt$$

$$= \int dt \dot{\theta} ml^2 \dot{\theta}$$
$$= \int dt ml^2 \dot{\theta}^2$$

$$\dot{\theta}(t) = A\omega_0 \sin(\omega_0 t)$$

$$= \int_{\text{cycle}} dt ml^2 A^2 \omega_0^2 \sin^2(\omega_0 t)$$

$$= \frac{1}{2} ml^2 A^2 \omega_0^2 (T) \quad T = 2\pi/\omega$$

$$= \pi ml^2 A^2 \omega_0$$

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$$\pi m l^2 A^2 \omega_0 = \text{const}$$

For small oscillations $\omega_0 = \sqrt{g/l}$

$$l^2 A^2 l^{-1/2} = \text{const}$$

$$A^2 l^{3/2} = \text{const}$$

$$A^2 \propto l^{-3/2}$$

$$A \propto l^{-3/4}$$

Solution for problem 2

Calculate the angular momentum stored in the fields before the current starts to flow; this will be the mechanical angular momentum after the capacitor is drained.

$$\vec{L}_{\text{em}} = \mu_o \epsilon_o \int \vec{r} \times \vec{S} dV = \epsilon_o \int \vec{r} \times \vec{E} \times \vec{B} dV$$

The electric field between the plates is:

$$\vec{E} = \frac{Q}{2\pi L \epsilon_o r} \hat{r}$$

And the Poynting vector in between the plates is then ($B = B\hat{z}$):

$$\vec{S} = -\frac{QB}{2\pi L \mu_o \epsilon_o r} \hat{\phi}$$

And L_{em} becomes (After integrating):

$$\vec{L}_{\text{em}} = -\frac{QB}{2\pi L} \int_a^b 2\pi r dr L = -\frac{QB}{2} (b^2 - a^2) \hat{z}$$

The final angular velocity of the capacitor is just this divided by the moment of inertia:

$$\vec{\omega} = \frac{\vec{L}_{\text{em}}}{I} = -\frac{QB}{2I} (b^2 - a^2) \hat{z}$$

Problem 3 Solutions:

a) At thermal equilibrium the temperatures on each side are equal: $T_{\text{He}} = T_{\text{SG}}$

helium: Since ideal gas:

$$U_{\text{He}} = \frac{3}{2} N k T_{\text{He}} \quad k = k_{\text{Boltzmann}}$$

$$T_{\text{He}} = \frac{2 U_{\text{He}}}{3 N k}$$

Supergas: $S_{\text{SG}} = k \ln \Omega$

$$S_{\text{SG}} = k \left[\ln U_{\text{SG}}^{10N} + \ln \left(\begin{smallmatrix} \text{other} \\ \text{stuff} \\ f(N, V) \end{smallmatrix} \right) \right]$$

$$\frac{1}{T_{\text{SG}}} = \left(\frac{\partial S_{\text{SG}}}{\partial U_{\text{SG}}} \right) \Big|_{V, N}$$

$$\frac{1}{T_{\text{SG}}} = 10 N k \left(\frac{1}{U_{\text{SG}}} \right)$$

$$T_{\text{SG}} = \frac{U_{\text{SG}}}{10 N k}$$

Using $T_{\text{He}} = T_{\text{SG}}$:

$$\frac{U_{\text{He}}}{3/2} = \frac{U_{\text{SG}}}{10} \Rightarrow \boxed{\frac{U_{\text{SG}}}{U_{\text{He}}} = \frac{20}{3}}$$

b) mechanical equilibrium \Rightarrow Final pressures are equal

helium:

$$P_{\text{He}} V_{\text{He}} = NkT \quad (\text{ideal gas})$$
$$P_{\text{He}} = \frac{NkT}{V_{\text{He}}}$$

Supergas:

$$P_{\text{SG}} = T \left(\frac{\partial S_{\text{SG}}}{\partial V_{\text{SG}}} \right) \bigg|_{U, N}$$
$$= T \left(\frac{k(10N)}{V_{\text{SG}}} \right)$$

$$P_{\text{He}} = P_{\text{SG}}$$

$$\frac{NkT}{V_{\text{He}}} = \frac{10NkT}{V_{\text{SG}}}$$

$$\boxed{\frac{V_{\text{SG}}}{V_{\text{He}}} = 10}$$

Solution to Problem #4

$$B = \frac{\mu_0 I}{4\pi} \int d\vec{\ell} \times \frac{(\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}$$

$$\vec{r} = 0$$

$$|\vec{r} - \vec{r}'|^2 = a^2 \phi'^2$$

$$\vec{r}' = \rho' \hat{\rho}' = a \phi' \hat{\rho}'$$



$$d\vec{\ell} = d\rho' \hat{\rho}' + \rho' d\phi' \hat{\phi}' =$$

$$= a d\phi' \hat{\rho}' + a \phi' d\phi' \hat{\phi}'$$

$$d\vec{\ell} \times \vec{r}' = \cancel{a^2 \phi' d\phi' \hat{\rho}' \times \hat{\rho}'} + a^2 \phi'^2 d\phi' \hat{\phi}' \times \hat{\rho}' =$$

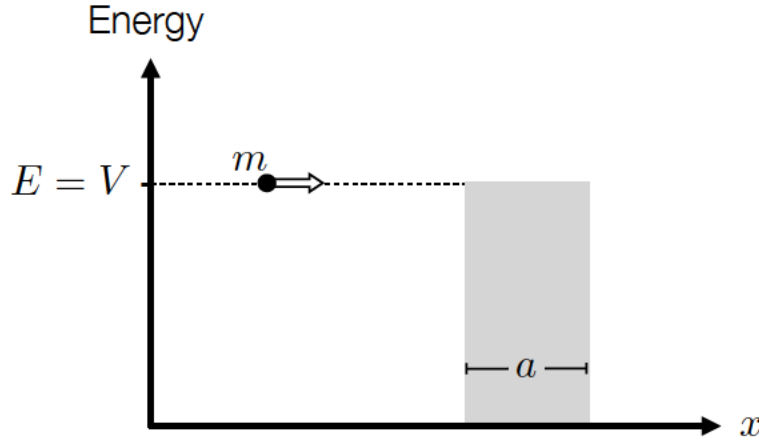
$$= -a^2 \phi'^2 d\phi' \hat{z}$$

$$B = \frac{\mu_0 I}{4\pi} \hat{z} \int_{2\pi}^{2\pi N} \frac{a^2 \phi'^2 d\phi'}{a^3 \phi'^3} = \frac{\mu_0 I}{4\pi a} \hat{z} \int_{2\pi}^{2\pi N} \frac{d\phi'}{\phi'} =$$

$$B = \frac{\mu_0 I}{4\pi a} \ln N \hat{z}$$

Solution: QM Problem 5

Starting with the 1D time-independent Schroedinger equation, work out the fraction of incident particles transmitted through a rectangular 1D potential barrier in the case that the energy E of the incident particles is equal to the barrier height V . Let the particles have mass m and let the barrier width be a .



Solution

We require solutions to the time-independent Schrodinger equation in three regions: (I) to the left of the barrier, (II) in the barrier where $E = V$, and (III) to the right of the barrier. The wavefunctions in these regions are

$$\psi_I = Ae^{ikx} + Be^{-ikx}, \quad \psi_{II} = Cx + D, \quad \psi_{III} = Fe^{ikx} \quad (1)$$

where $k = \sqrt{2mE/\hbar^2}$. These wavefunctions and their derivatives are continuous at the boundaries of the barrier. Let the left boundary be at $x = 0$ and the right boundary be at $x = a$. Matching gives the equations

$$A + B = D, \quad ik(A - B) = C, \quad Ca + D = Fe^{ika}, \quad C = ikFe^{ika} \quad (2)$$

We want to find F/A . Using these equations we get,

$$\frac{F}{A} = \frac{e^{-ika}}{1 - ika/2} \quad (3)$$

The transmission coefficient is then

$$T = \left| \frac{F}{A} \right|^2 = \frac{1}{1 + \frac{mEa^2}{2\hbar^2}} \quad (4)$$

Solution: Problem 6

(Source: This is almost verbatim from the discussion in Ch. 1 of Sakurai that concludes with his statement, “Here lies the heart of quantum mechanics.”)

- (a) Operating on the left with any arbitrary B eigenstate $\langle b''|$ gives

$$\begin{aligned}\langle b''|a'\rangle &= \langle b''|\sum_b c_b|b\rangle \\ &= c_{b''}\langle b''|b''\rangle \\ &= c_{b''}\end{aligned}$$

and so $c_b = \langle b|a'\rangle$. The same procedure applied to the second case shows that $c_c = \langle c|b'\rangle$.

- (b) The probability of outcome $|b'\rangle$ given $|a'\rangle$ is $P(b'|a') = |c_{b'}|^2$. The probability of outcome $|c'\rangle$ given $|b'\rangle$ is $P(c'|b') = |c_{c'}|^2$. Therefore, we have $P(c'|a' \& b') = |c_{b'}|^2 |c_{c'}|^2$.

- (c)

$$\begin{aligned}\sum_{b'} P(c'|a' \& b') &= \sum_{b'} |c_{b'}|^2 |c_{c'}|^2 \\ &= \sum_{b'} \langle a'|b'\rangle \langle b'|c'\rangle \langle c'|b'\rangle \langle b'|a'\rangle\end{aligned}\tag{3}$$

- (d) $P(c'|a') = |\langle a'|c'\rangle|^2 = \langle a'|c'\rangle \langle c'|a'\rangle$. To compare this to Eq. (3), we can write $\langle a'|$ in the B basis,

$$\langle a'| = \sum_b \langle a'|b\rangle \langle b|$$

which gives us

$$\begin{aligned}P(c'|a') &= \left| \sum_b \langle a'|b\rangle \langle b|c'\rangle \right|^2 \\ &= \sum_{b,b''} \langle a'|b\rangle \langle b|c'\rangle \langle c'|b''\rangle \langle b''|a'\rangle.\end{aligned}\tag{4}$$

Comparing Eq. (3) to Eq. (4), we see that these will not have to be the same probability.

There are other ways to arrive at this answer as well. For instance, if $|a'\rangle$ and $|c'\rangle$ are orthogonal, clearly $P(c'|a') = 0$, but there may be B eigenstates that are not orthogonal to either one, in which case result (3) is nonzero.