

(a) Percolation threshold is a mathematical concept related to percolation theory, which is the formation of long-range connectivity in random systems. Below the threshold a giant connected component does not exist;

Conductivity near the percolation threshold is a mixture between a dielectric and a metallic component, the conductivity  $\sigma$  and the dielectric constant  $\epsilon$  of this mixture show a critical behaviour if the fraction of the metallic component reaches the percolation threshold.

$$\xi(p) \propto (p_c - p)^{-\nu}$$

$$P(p) \propto (p - p_c)^{\beta}$$

A one-dimensional chain state of vortex matters. The zero-dimensional analogs of the two-dimensional states that occur at a crystal surface were observed at the ends of one-dimensional chain.

(b) In an infinite one-dimensional chain, the probability  $n_s$  that a given site is the left end of a cluster of length precisely  $s$ .

$$P(n \text{ arbitrary sites are occupied}) = p^n.$$

The probability that a fixed site is to the left end of an  $s$ -cluster is therefore  $(1-p)^2 p^s$ .

Each site has an equal probability of being the left end of an  $s$ -cluster. Thus in a system of size  $L$ , the total number of  $s$ -clusters is, ignoring end effects,

$$L(1-p)^2 p^s.$$

$n_s =$  number of  $s$ -clusters per site.

$$\text{then, } n_s = p^s (1-p)^2.$$

We also have

$$P(\text{arbitrary site is part of } s\text{-cluster}) = n_s s,$$

larger than  $n_s$  by a factor of  $s$ , since there are  $s$  ways an arbitrary site can be situated within an  $s$ -cluster.

(c) Summing over all possible cluster sizes, we obtain the probability that an arbitrary site is a member of any cluster.

$$\sum_{s=1}^{\infty} n_s s = p \quad (p < p_c).$$

where we write  $p < p_c$  to eliminate the possibility of an infinite cluster at  $p_c$ .

It is instructive to obtain the previous result from the definition of  $n_s$

$$\begin{aligned} \sum_s n_s s &= \sum_s p^s (1-p)^2 s \\ &= (1-p)^2 \sum_s p \frac{d}{dp} p^s \\ &= (1-p)^2 p \frac{d}{dp} \sum_s p^s \\ &= (1-p)^2 p \frac{d}{dp} \left( \frac{p}{1-p} \right) \\ &= p \end{aligned}$$

(d) Another quantity of interest is

$$\omega_s = P(\text{arbitrary occupied site belongs to an } s\text{-cluster})$$

$$\begin{aligned} &= \frac{P(\text{site belongs to } s\text{-cluster})}{P(\text{site belongs to any cluster})} \\ &= \frac{n_s S}{\sum_s n_s S} \end{aligned}$$

Then, defining  $S$  = average cluster size,

$$S = \sum_s \omega_s S = \sum_s \left( \frac{n_s S^2}{\sum_s n_s S} \right)$$

$$\begin{aligned} \sum_s n_s S^2 &= (1-p)^2 \sum_s s^2 p^s \\ &= (1-p)^2 \left( p \frac{d}{dp} \right)^2 \sum_s p^s \end{aligned}$$

that  $\left( p \frac{d}{dp} \right)^2 = p \frac{d}{dp} + p^2 \frac{d^2}{dp^2}$  and using the previous trick to calculate the sums, we obtain

$$S = \frac{1+p}{1-p}$$

and note that  $S$  diverges at  $p = p_c = 1$ .